

Lecture 3

1 How to construct invariants for a given group action?

Let the action given by a single vector field \mathbf{X} in \mathbb{R}^m . If $\mathbf{X} \neq 0$ everywhere then there exists exactly $m - 1$ functionally independent invariants $\phi_1, \dots, \phi_{m-1}$ for this action.

If s vector fields are given which span s -dimensional orbits of a group action, then we expect to find $m - s$ functionally dependent invariants.

Examples:

- $SO(2)$ acting on \mathbb{R}^2 via $\mathbf{X} = x\partial_y - y\partial_x$. Then $x^2 + y^2$ is an invariant, and any other invariant is a function thereof.
- $SO(2)$ acting on \mathbb{R}^3 via $\mathbf{X} = x\partial_y - y\partial_x + (1 + z^2)\partial_z$. Then we have to solve

$$\frac{dx}{-y} = \frac{dy}{x} = \frac{dz}{1 + z^2}$$

Since $r^2 = x^2 + y^2$ is an invariant, then we can eliminate $x = \sqrt{r^2 - y^2}$ so

$$\frac{dy}{\sqrt{r^2 - y^2}} = \frac{dz}{1 + z^2}$$

which implies

$$\arcsin(y/r) = \arctan(z) + k$$

but $y/r = \sin \theta$ implies $x/y = \tan \theta$ so $\arcsin(y/r) = \arctan(y/x)$, hence

$$\arctan(z) - \arctan(y/x)$$

is a second invariant. So

$$\tan(\arctan(z) - \arctan(y/x)) = \frac{z - y/x}{1 + zy/x} = \frac{zx - y}{yz + x} := q \quad (1.1)$$

is a second invariant (functionally dependent on the former). It is defined on the domain $z \neq -x/y$. Alternatively:

$$\tilde{q} = \frac{r}{\sqrt{1 + q^2}} = \frac{x + yz}{\sqrt{1 + z^2}}$$

is an invariant defined everywhere.

- $\mathbf{X}_1 = -y\partial_x + x\partial_y$, $\mathbf{X}_2 = 2xz\partial_x + 2yz\partial_y + (z^2 + 1 - x^2 - y^2)\partial_z$. We have $[\mathbf{X}_1, \mathbf{X}_2] = 0$ so these generate an abelian group. We already know an invariant r of \mathbf{X}_1 . Is there another invariant? Yes, it is ∂_z . Since any invariant of $(\mathbf{X}_1, \mathbf{X}_2)$ is, in particular, an invariant of \mathbf{X}_1 . it must be a function $\phi(r, z)$. So, we may push \mathbf{X}_2 to the coordinates (r, z) , that is, consider $F(x, y, z) = (r, z)$ and let $\mathbf{Y} := F_*(\mathbf{X}_2)$. We calculate \mathbf{Y} easily via

$$\mathbf{Y} = \mathbf{X}_2(r)\partial_r + \mathbf{X}_2(z)\partial_z = 2rz\partial_r + (z^2 + 1 - r^2)\partial_z$$

The function we are looking satisfies

$$2rz \frac{\partial \phi}{\partial r} + (z^2 + 1 - r^2) \frac{\partial \phi}{\partial z} = 0$$

which verifies the characteristic equation

$$\frac{dr}{2rz} = \frac{dz}{z^2 + 1 - r^2}$$

The solution is

$$\phi = \frac{z^2 + r^2 + 1}{r} = \frac{x^2 + y^2 + z^2 + 1}{\sqrt{x^2 + y^2}}$$

2 Invariance of Differential equations: Overview

Definition 2.1. An ODE of order n is a relation

$$H(x, y, y', \dots, y_n) = 0 \quad (2.1)$$

where $H(x_1, \dots, x_{n+1}) : \mathbb{R}^{n+1} \rightarrow \mathbb{R}$.

A point transformation is an action $\Psi := (\psi_1, \psi_2)$ of a Lie group G acting on the space \mathbb{R}^2 identified with (x, y) : $\Psi(g, x, y) = (\psi_1(g, x, y), \psi_2(g, x, y)) := (x_g, y_g)$.

The ODE (2.1) is invariant under the point transformation iff for any solution $y = y(x)$ of the ODE, the graph $(x, y(x))$ is transformed under Ψ to a graph of another solution of the same ODE: $(\psi_1(g, x, y(x)), \psi_2(g, x, y(x)))$ is a graph of a solution $y_{(g)} = y_{(g)}(x)$ for any $g \in G$.

Example: The action of $SO(2)$ on \mathbb{R}^2 via

$$\Psi(t, x, y) = (x \cos(t) - y \sin(t), x \sin(t) + y \cos(t))$$

preserves the linear functions $y = ax + b$. Indeed,

$$(x_{(t)}, y_{(t)}) = (x \cos(t) - (ax + b) \sin(t), x \sin(t) + (ax + b) \cos(t))$$

and we eliminate

$$x = \frac{x_{(t)} + b \sin(t)}{\cos(t) - a \sin(t)}$$

and

$$y_{(t)} = x \sin(t) + (ax + b) \cos(t) = \frac{\sin(t) + a \cos(t)}{\cos(t) - a \sin(t)} x_{(t)} + \frac{b}{\cos(t) - a \sin(t)}$$

Hence: The equation $y'' = 0$ is invariant under the $SO(2)$ action.

If we replace the slope $a = y'$ by y_1 , we see that the action $SO(2)$ is extended to an action on \mathbb{R}^3 parameterized by x, y, y_1 as:

$$(x, y, y_1) \rightarrow \left(x \cos(t) - y \sin(t), x \sin(t) + y \cos(t), \frac{\sin(t) + y_1 \cos(t)}{\cos(t) - y_1 \sin(t)} \right) \quad (2.2)$$

which is induced by the vextorfield

$$\mathbf{X} = -y\partial_x + x\partial_y + (1 + (y_1)^2)\partial_{y_1}$$

Let now $y = y(x)$ be any function. Suppose $y = ax + b$ is the tangent to its graph at (x_0, y_0) . Then

$$y_{1,(t)} = a_{(t)} := \frac{\sin(t) + a \cos(t)}{\cos(t) - a \sin(t)}$$

is the slop to the tangent of the transformed graph at $(x_{(t)}, y_{(t)})$. In particular, we obtained that the derivative y' is transformed as y_1 in the extended action.

We shall see later that the invariants of this action on \mathbb{R}^3 , namely

$$r^2 = x^2 + y^2 \quad , \quad q = \frac{y'x - y}{yy' + x}$$

(see (1.1)) are constants for any solution of $SO(2)$ invariant equation:

$$r^2(x, y(x)) = C_1, \quad q(x, y(x), y'(x)) = C_2 \quad ,$$

and that any such equation of first order is given by a functional relation between r^2 and q :

$$H(x, y, y') = F(r^2, q) = 0 \quad .$$

3 Prolongation of vector fields

Motivated by the extension of functional to differential invariance we wish to extend the graph of a function $y = f(x)$ in \mathbb{R}^2 to a graph of a *prolonged* function $(y, y') = Pr^{(1)}f(x) := (f(x), f'(x))$ in \mathbb{R}^3 . More generally,

$$Y^{(n)} = (y, y', y'', \dots, y_n) = (f(x), f'(x), \dots, f^{(n)}(x))$$

What is the image of $Y^{(n)} = Pr^{(n)}f(x)$ under the flow of a vector field $\mathbf{X} = \xi\partial_x + \eta\partial_y$?

$$\tilde{y}' = \frac{d\tilde{y}}{d\tilde{x}} \quad ; \quad \widetilde{y_{k+1}} = \frac{d\tilde{y}_k}{d\tilde{x}}$$

and so on, where

$$\tilde{x} = x + \varepsilon\xi(x, y) + \dots \quad , \quad \tilde{y} = y + \varepsilon\eta(x, y) + \dots \quad ; \quad \tilde{y}_1 = y_1 + \varepsilon\eta^{(1)}(x, y, y_1) + \dots$$

Here we make, for any $h = h(x, y, y_1, \dots$

$$dh = (h_x + h_y y_1 + \dots h_{y_n} y_{n+1}) dx$$

implies

$$\begin{aligned} d\tilde{y} &= dy + \varepsilon d\eta \quad , \quad d\tilde{x} = dx + \varepsilon d\xi \\ \tilde{y}_1 &= \frac{d\tilde{y}}{d\tilde{x}} = \frac{dy}{dx} + \varepsilon \left(\frac{d\eta}{dx} - y_1 \frac{d\xi}{dx} \right) := y_1 + \varepsilon\eta^{(1)} + \end{aligned}$$

as well as

$$\widetilde{y_n} = \frac{d\tilde{y}_{n-1}}{d\tilde{x}} = \frac{dy_{n-1}}{dx} + \varepsilon \left(\frac{d\eta^{(n-1)}}{dx} - y_n \frac{d\xi}{dx} \right) := y_n + \varepsilon\eta^{(n)} +$$

Definition 3.1. The prolongation of a vector field $\mathbf{X} = \xi\partial_x + \eta\partial_y$ to order n is given by

$$Pr^{(n)}\mathbf{X} := \xi\partial_x + \eta\partial_y + \eta^{(1)}\partial_{y_1} + \dots + \eta^{(n)}\partial_{y_n}$$

where

$$\eta^{(0)} \equiv \eta, \quad \eta^{(k)} = \frac{d\eta^{(k-1)}}{dx} - y_k \frac{d\xi}{dx}$$

We shall sometimes write \mathbf{X} instead of $Pr^{(n)}\mathbf{X}$, when no confusion is expected.

Examples:

$$(A) -y\partial_x + x\partial_y + (1 + (y_1)^2)\partial_{y_1} + 3y_1y_2\partial_{y_2} + (3(y_2)^2 + 4y_1y_3)\partial_{y_3} + \dots$$

$$(B) x\partial_x + y\partial_y - y_2\partial_{y_2} - 2y_3\partial_{y_3} \dots$$

$$(C) \partial_x, \quad x\partial_x - y_1\partial_{y_1} - 2y_2\partial_{y_2} - 3y_3\partial_{y_3} + \dots, \quad x^2\partial_x - 2xy_1\partial_{y_1} - (2y_1 + 4xy_2)\partial_{y_2} - 6(xy_3 - y_2)\partial_{y_3} + \dots$$

Theorem 1. The prolongation of vector fields is compatible with the Lie product, i.e

$$[Pr^{(n)}(\mathbf{X}_1), Pr^{(n)}(\mathbf{X}_2)] = Pr^{(n)}[\mathbf{X}_1, \mathbf{X}_2] .$$

In particular, it follows that the prolongation of a Lie algebra is an isomorphic Lie algebra.

Definition 3.2. The n - prolongation of action $\Psi : G \times \mathbb{R}^2 \rightarrow \mathbb{R}^2$ is denoted by $Pr^{(n)}\Psi : G \times \mathbb{R}^{2+n} \rightarrow \mathbb{R}^{2+n}$. The invariants vectorfields are the prolonged Lie algebra corresponding to Ψ .

Remark 3.1. Note that, in general, the prolonged action is only local (even if the original action is global- see (2.2)).